

5.1.1 Similar Formulas

$$\cos(\theta) = \frac{e^{i\theta} + e^{-i\theta}}{2} \quad (5.7)$$

$$\sin(\theta) = \frac{e^{i\theta} - e^{-i\theta}}{2i} \quad (5.8)$$

$$\cos(z) = \frac{e^{iz} + e^{-iz}}{2} \quad (5.9)$$

$$\sin(z) = \frac{e^{iz} - e^{-iz}}{2} \quad (5.10)$$

$$\cosh z = \frac{e^z + e^{-z}}{2} \quad (5.11)$$

$$\sinh z = \frac{e^z - e^{-z}}{2} \quad (5.12)$$

$$\cosh iz = \cos z \quad (5.13)$$

$$\sinh iz = i \sin z \quad (5.14)$$

5.1.2 De Moivre's Theorem

$$(\cos \theta + i \sin \theta)^n = \cos n\theta + i \sin n\theta \quad (5.15)$$

From this it can be shown with $z = e^{i\theta}$

$$z^n + \frac{1}{z^n} = 2 \cos n\theta \quad z^n - \frac{1}{z^n} = 2i \sin n\theta \quad (5.16)$$

5.1.3 Roots of unity

$$z^n = 1 = e^{i2k\pi} \Rightarrow z = e^{i2k\pi/n} \quad (5.17)$$

where k is an integer and will take values $0, 1, 2, \dots, n-1$ to give the n distinct roots

5.1.4 Logarithms

with $z = re^{i(\theta+2n\pi)}$

$$Ln z = \ln r + i(\theta + 2n\pi) \quad (5.18)$$

Restricting to principal value by constraining the argument of z to lie between $-\pi$ to π we get singlevalued $Ln(z)$

The definition of complex number raised to a complex power is in terms of already defined complex functions:

$$t^z = e^{z \ln(t)} \quad (5.19)$$

5.2 Complex Functions

$$\lim_{z \rightarrow z_0} f(z) = \infty \text{ if and only if } \lim_{z \rightarrow z_0} \frac{1}{f(z)} = 0 \quad (5.20)$$

$$\lim_{z \rightarrow \infty} f(z) = w_0 \text{ if and only if } \lim_{z \rightarrow 0} f\left(\frac{1}{z}\right) = w_0 \quad (5.21)$$

$$\lim_{z \rightarrow \infty} f(z) = \infty \text{ if and only if } \lim_{z \rightarrow 0} \frac{1}{f(1/z)} = 0 \quad (5.22)$$

Since x and y are related to z and its complex conjugate z^* by

$$x = \frac{1}{2}(z + z^*) \quad \text{and} \quad y = \frac{1}{2i}(z - z^*), \quad (24.6)$$

we may formally regard any function $f = u + iv$ as a function of z and z^* , rather than x and y . If we do this and examine $\partial f / \partial z^*$ we obtain

$$\begin{aligned} \frac{\partial f}{\partial z^*} &= \frac{\partial f}{\partial x} \frac{\partial x}{\partial z^*} + \frac{\partial f}{\partial y} \frac{\partial y}{\partial z^*} \\ &= \left(\frac{\partial u}{\partial x} + i \frac{\partial v}{\partial x} \right) \left(\frac{1}{2} \right) + \left(\frac{\partial u}{\partial y} + i \frac{\partial v}{\partial y} \right) \left(-\frac{1}{2i} \right) \\ &= \frac{1}{2} \left(\frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \right) + \frac{i}{2} \left(\frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right). \end{aligned} \quad (24.7)$$

Now, if f is analytic then the Cauchy-Riemann relations (24.5) must be satisfied, and these immediately give that $\partial f / \partial z^*$ is identically zero. Thus we conclude that if f is analytic then f cannot be a function of z^* and any expression representing an analytic function of z can contain x and y only in the combination $x + iy$, *not* in the combination $x - iy$.

5.2.1 Cauchy-Reimann Equations

Suppose that for a complex-valued function

$$f(z) = u(x, y) + iv(x, y) \quad (5.23)$$

$f'(z)$ exists at a point $z_0 = (x_0, y_0)$. Then the first-order partial derivatives of u and v must exist at (x_0, y_0) and they must satisfy the Cauchy-Reimann Equations:

$$u_x = v_y, \quad u_y = -v_x \quad (5.24)$$

And also

$$f'(z_0) = u_x + iv_x \quad (5.25)$$

- Differentiation of Cauchy equations provides us with the fact that both u and v individually satisfy Laplace equations in 2D

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0 \quad \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = 0 \quad (5.26)$$

- The family of 2D curves (in xy plane) $u(x, y) = \text{constant}$ and $v(x, y) = \text{constant}$ intersect at right angles to each other.
- For multivalued complex function $f(z)$, like $\exp(z)$, $\text{Ln} z$, $z^{1/2}$ etc. we can make them single valued and then use the analytic function analysis on it by concept of branch points and cuts
- Branch point is a point in Argand plane such that if z is varied in a closed loop enclosing the branch point, $f(z)$ doesn't return to its original value, although z does since $\theta \rightarrow \theta + 2\pi$ doesn't affect z .
- Branch cuts are lines in the Argand diagram, finite or infinite that prevents us to ever make a closed loop enclosing a branch point.
- So, if we don't cross the branch cut, $f(z)$ remains single-valued.

5.2.2 Singularities

- Isolated singularity: when $f(z)$ isn't analytic at z_0 but analytic at all points in the neighbourhood.
- Pole: most imp isolated singularity
- If

$$f(z) = \frac{g(z)}{(z - z_0)^n} \quad (5.27)$$

such that $g(z_0) \neq 0$ and is analytic in neighbourhood of z_0 then it's a pole of order n

$$\nabla \times \mathbf{A} = \frac{1}{r} \begin{vmatrix} \hat{\mathbf{r}} & r\hat{\boldsymbol{\theta}} & \hat{\mathbf{z}} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial z} \\ A_1 & rA_2 & A_3 \end{vmatrix} \quad (6.11)$$

Spherical

$$\nabla f = \frac{\partial f}{\partial r} \hat{\mathbf{r}} + \frac{1}{r \sin \phi} \frac{\partial f}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r} \frac{\partial f}{\partial z} \hat{\boldsymbol{\phi}} \quad (6.12)$$

$$\nabla \cdot \mathbf{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_1) + \frac{1}{r \sin \phi} \frac{\partial A_2}{\partial \theta} + \frac{1}{r \sin \phi} \frac{\partial}{\partial \phi} (\sin \phi A_3) \quad (6.13)$$

$$\nabla \times \mathbf{A} = \frac{1}{r^2 \sin \phi} \begin{vmatrix} \hat{\mathbf{r}} & r \sin \phi \hat{\boldsymbol{\theta}} & r \hat{\boldsymbol{\phi}} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ A_1 & r \sin \phi A_2 & r A_3 \end{vmatrix} \quad (6.14)$$

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial^2}{\partial \phi^2} \right) \quad (6.15)$$

6.2.5 Fundamental theorems

The line integral of any gradient field doesn't depend on path but just on the endpoints. The little changes df add up (the integral on LHS) to give $f(b) - f(a)$

$$\int_a^b (\nabla f) \cdot d\mathbf{l} = f(b) - f(a) \quad (6.16)$$

The volume integral of a divergence over any given volumen is equal to the closed surface integral or the flux of the original vector field over the surface that bounds that volume

$$\int_V (\nabla \cdot \mathbf{A}) dV = \int_S \mathbf{A} \cdot d\mathbf{a} \quad (6.17)$$

The surface integral of curl of a vector field is equal to the closed line integral of that field over the boundary of the surface. Since for a given boundary there can be infinite number of surfaces, it shows that surface integral of curls is only dependent on the boundary of the surface.

$$\int_S (\nabla \times \mathbf{A}) \cdot d\mathbf{a} = \int \mathbf{A} \cdot d\mathbf{l} \quad (6.18)$$

These operators are usually applied on scalar (gradient) and vector (divergence, curl) fields. Boldface notation is used for vector fields. f and g are scalar fields and k is just a scalar number

6.2.6 First derivatives

- Sum rules

$$\nabla(f + g) = \nabla f + \nabla g \quad (6.19)$$

$$\nabla \cdot (\mathbf{A} + \mathbf{B}) = (\nabla \cdot \mathbf{A}) + (\nabla \cdot \mathbf{B}) \quad (6.20)$$

$$\nabla \times (\mathbf{A} + \mathbf{B}) = (\nabla \times \mathbf{A}) + (\nabla \times \mathbf{B}) \quad (6.21)$$

- Multiplying by a constant rules

$$\nabla(kf) = k\nabla f \quad (6.22)$$

$$\nabla \cdot (k\mathbf{A}) = k(\nabla \cdot \mathbf{A}) \quad (6.23)$$

$$\nabla \times (k\mathbf{A}) = k(\nabla \times \mathbf{A}) \quad (6.24)$$

Chapter 7

Coordinate Systems

7.1 Cartesian coordinates

Straightforward and easy to understand formulas for all things.

- Point P(x,y,z)

where x, y and z denotes the distances along the x,y and z axes respectively.

The position vector of any point P is given by $OP = x\hat{x} + y\hat{y} + z\hat{z}$ where the unit vectors are $\hat{x}$, $\hat{y}$ and $\hat{z}$ that are fixed in magnitude(= 1) and direction also. So in any problem with changing position/displacement vector, their time derivative must be zero

- Distance Formula

Between two points $P(x_1, y_1, z_1)$ and $Q(x_2, y_2, z_2)$ is:

$$PQ = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2} \quad (7.1)$$

- Equation of plane with unit normal vector $\hat{n}$ and a given point in plane with position vector $\mathbf{r}'$:

$$(\mathbf{r} - \mathbf{r}') \cdot \hat{n} = 0 \quad (7.2)$$

- In component form:

$$px + qy + sz = d \quad \hat{n} = (p, q, s) \quad \mathbf{r} \cdot \hat{n} = d \quad (7.3)$$

where d = perpendicular distance of plane from origin

- In intercept form:

$$\frac{x}{a} + \frac{y}{b} + \frac{z}{c} = 1 \quad (7.4)$$

- Relation b/w unit normal and intercepts:

$$\hat{n} \propto \frac{\hat{i}}{a} + \frac{\hat{j}}{b} + \frac{\hat{k}}{c} \quad (7.5)$$

- Angle between two planes: $l_1x + m_1y + n_1z = d_1$ and $l_2x + m_2y + n_2z = d_2$ is:

$$\cos \theta = \frac{l_1l_2 + m_1m_2 + n_1n_2}{\sqrt{l_1^2 + m_1^2 + n_1^2} \sqrt{l_2^2 + m_2^2 + n_2^2}} \quad (7.6)$$

- Laplacian

$$\nabla \cdot (\nabla f) = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} \quad (7.7)$$

- Laplacian of a vector

$$\nabla^2 \mathbf{v} \equiv (\nabla^2 v_x)\hat{x} + (\nabla^2 v_y)\hat{y} + (\nabla^2 v_z)\hat{z} \quad (7.8)$$

3D spherical polar

Line element:

$$ds^2 = dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \quad (7.21)$$

Volume element:

$$dV = r^2 \sin \theta d\phi d\theta dr \quad (7.22)$$

Cylindrical

Line element:

$$ds^2 = d\rho^2 + \rho^2 d\phi^2 + dz^2 \quad (7.23)$$

Volume element:

$$dV = \rho d\phi d\rho dz \quad (7.24)$$

7.4 Coordinate transformations and unit vectors

SPHERICAL AND CYLINDRICAL COORDINATES

Spherical

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases} \quad \begin{cases} \hat{\mathbf{x}} = \sin \theta \cos \phi \hat{\mathbf{r}} + \cos \theta \cos \phi \hat{\boldsymbol{\theta}} - \sin \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{y}} = \sin \theta \sin \phi \hat{\mathbf{r}} + \cos \theta \sin \phi \hat{\boldsymbol{\theta}} + \cos \phi \hat{\boldsymbol{\phi}} \\ \hat{\mathbf{z}} = \cos \theta \hat{\mathbf{r}} - \sin \theta \hat{\boldsymbol{\theta}} \end{cases}$$

$$\begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \theta = \tan^{-1} \left(\sqrt{x^2 + y^2} / z \right) \\ \phi = \tan^{-1}(y/x) \end{cases} \quad \begin{cases} \hat{\mathbf{r}} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}} \\ \hat{\boldsymbol{\theta}} = \cos \theta \cos \phi \hat{\mathbf{x}} + \cos \theta \sin \phi \hat{\mathbf{y}} - \sin \theta \hat{\mathbf{z}} \\ \hat{\boldsymbol{\phi}} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \end{cases}$$

Cylindrical

$$\begin{cases} x = s \cos \phi \\ y = s \sin \phi \\ z = z \end{cases} \quad \begin{cases} \hat{\mathbf{s}} = \cos \phi \hat{\mathbf{x}} + \sin \phi \hat{\mathbf{y}} \\ \hat{\boldsymbol{\phi}} = -\sin \phi \hat{\mathbf{x}} + \cos \phi \hat{\mathbf{y}} \\ \hat{\mathbf{z}} = \hat{\mathbf{z}} \end{cases}$$

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12.2 Fourier integral and transforms

Fourier cosine integral which applies to an even function $f(x)$

$$A(w) = \frac{2}{\pi} \int_0^{\infty} f(x) \cos wx \, dx \quad f(x) = \int_0^{\infty} A(w) \cos wx \, dw \quad (12.6)$$

Fourier sine integral which applies to an odd function $f(x)$

$$B(w) = \frac{2}{\pi} \int_0^{\infty} f(x) \sin wx \, dx \quad f(x) = \int_0^{\infty} B(w) \sin wx \, dw \quad (12.7)$$

For a general function $f(x)$, we have representation in terms of its fourier integral. Here the series of eq. 11.1 transforms to an integral as we take $L \rightarrow \infty$

$$f(x) = \int_0^{\infty} [A(w) \cos wx + B(w) \sin wx] dw \quad (12.8)$$

$$A(w) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \cos wv \, dv, \quad B(w) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(v) \sin wv \, dv \quad (12.9)$$

Writing in complex and doing few tricks we get:

$$\hat{f}(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx \quad (12.10)$$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(k) e^{ikx} dk \quad (12.11)$$

12.2.1 Fourier sine and cosine transforms

$$\hat{f}_s(k) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(x) \sin kx \, dx \quad (12.12)$$

$$\hat{f}_c(k) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{f}_s(k) \sin kx \, dk \quad (12.13)$$

$$\hat{f}_c(k) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(x) \cos kx \, dx \quad (12.14)$$

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{f}_c(k) \cos kx \, dk \quad (12.15)$$

Properties

$$\mathcal{F}_s\{f'(x)\} = -k\mathcal{F}_c\{f(x)\} \quad (12.16)$$

$$\mathcal{F}_c\{f'(x)\} = k\mathcal{F}_s\{f(x)\} - \sqrt{\frac{2}{\pi}} f(0) \quad (12.17)$$

$$\mathcal{F}_c\{f''(x)\} = -k^2\mathcal{F}_c\{f(x)\} - \sqrt{\frac{2}{\pi}} f'(0) \quad (12.18)$$

$$\mathcal{F}_s\{f''(x)\} = -k^2\mathcal{F}_s\{f(x)\} + \sqrt{\frac{2}{\pi}} k f(0) \quad (12.19)$$

$$\frac{d}{dk} \mathcal{F}_s(k) = \mathcal{F}_c\{xf(x)\} \quad (12.20)$$

$$\frac{d}{dk} \mathcal{F}_c(k) = -\mathcal{F}_s\{xf(x)\} \quad (12.21)$$

13.1.3 Poisson distribution

This distribution is actually a discrete distribution cause the random variable is discrete. But that random variable is spread over a continuum. Like the number of calls that we will receive over some particular time interval.

$$P_x(t) = \frac{(\lambda t)^x}{x!} e^{-\lambda t} \quad (13.22)$$

This is the probability of receiving x number of calls in a time period of t given that λ is the average number of calls per unit time.

Mean and variance are both λ

13.1.4 Gaussian distribution

Many random variables occurring in physical sciences and others follow this distribution exactly or approximately.

The probability density function is:

$$f(x) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2\right] \quad (13.23)$$

Mean and variance can be changed by shift of origin and scaling. Then with standard variable $Z = (X - \mu)/\sigma$

$$\phi(z) = \frac{1}{\sqrt{2\pi}} \exp\left(-z^2/2\right) \quad (13.24)$$

The cumulative probability function cant be solved analytically but is tabulated as $\Phi(z)$ for different z values.

$$F(x) = \Phi\left(\frac{x-\mu}{\sigma}\right) \quad (13.25)$$

$$Pr(a < X \leq b) = \Phi\left(\frac{b-\mu}{\sigma}\right) - \Phi\left(\frac{a-\mu}{\sigma}\right) \quad (13.26)$$

13.2 Random walk

- Simplest random walk: Walk through steps by flipping a coin. If head comes move forward, tail comes move backwards.
- Parameters: N = total number of steps taken. l = length of each step. n = the integer (if $l = 1$) where you landed at the end.
- Pattern on analysis probabilities of different possible ns when N changes.

n	-5	-4	-3	-2	-1	0	1	2	3	4	5
$f_0(n)$						1					
$2f_1(n)$					1		1				
$2^2f_2(n)$				1		2		1			
$2^3f_3(n)$			1		3		3		1		
$2^4f_4(n)$		1		4		6		4		1	
$2^5f_5(n)$	1		5		10		10		5		1

- Thus, the coefficients are binomial
- If number of heads = $n_{forward}$ and tails = $n_{backward}$ get fixed then final position is fixed $n = n_{forward} - n_{backward}$.